

# Oscillation modes, tides, and resonances of compact objects (2nd part)

Sayan Chakrabarti<sup>1</sup>    T rence Delsate<sup>2</sup>    Jan Steinhoff<sup>1,3</sup>



<sup>1</sup>Centro Multidisciplinar de Astrof sica (CENTRA)  
Instituto Superior T cnico (IST), Lisbon

<sup>2</sup>UMons  
Mons, Belgium

<sup>3</sup>ZARM  
University of Bremen

Mons Meeting 2013, PandA Doctoral School, July 15th, 2013, Mons

Supported by **DFG** through STE 2017/1-1 and STE 2017/2-1 “Resonances of quasinormal modes and orbital motion in general relativistic compact binaries” and by **FCT** (Portugal)

# Newtonian case

- Amplitudes:

$$\ddot{A}_{nlm} + \omega_{nl}^2 A_{nlm} = f_{nlm}$$

- Overlap integrals  $I_{nl}$ :

$$f_{nlm} = -\frac{1}{l!} I_{nl} \Phi_{\text{ext}}^{lm}, \quad q^{lm} = \sum_n I_{nl} A_{nlm}$$

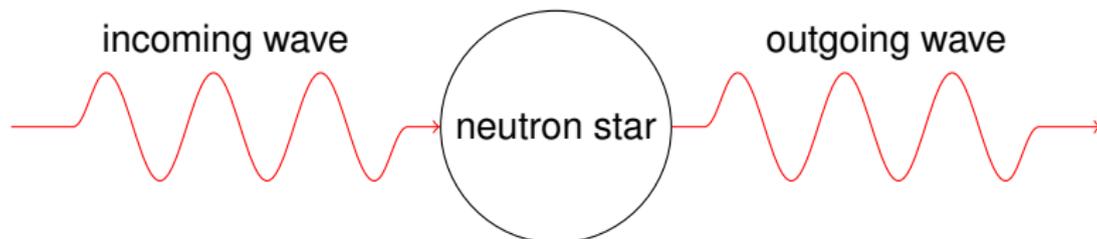
- Linear response in frequency domain:

$$\tilde{q}^{lm}(\omega) = -\frac{1}{l!} \tilde{F}_l(\omega) \tilde{\Phi}_{\text{ext}}^{lm}(\omega)$$

- Newtonian case: **response is sum of harmonic oscillators**

$$\tilde{F}_l(\omega) = \sum_n \frac{I_{nl}^2}{\omega_n^2 - \omega^2} \quad [\text{Chakrabarti, Delsate, Steinhoff, arXiv:1306.5820}]$$

- From now on:  $l = 2$**



- Inhom. Regge-Wheeler eq. with **effective  $\delta$ -source  $S$**  representing a NS

$$\frac{d^2 X}{dr_*^2} + \left[ \left(1 - \frac{2M}{r}\right) \frac{l(l+1) - \frac{6M}{r}}{r^2} + \omega^2 \right] X = S$$

- Analytic solutions for hom. eq. [Mano, Suzuki, Takasugi, PTP **96** (1996) 549]

$$X_{UV}^\nu = e^{-i\omega r} (\omega r)^\nu \left(1 - \frac{2M}{r}\right)^{-i2M\omega} \sum_{n=-\infty}^{\infty} \cdots \times \left[\frac{r}{2M}\right]^n {}_2F_1(\dots; 2M/r)$$

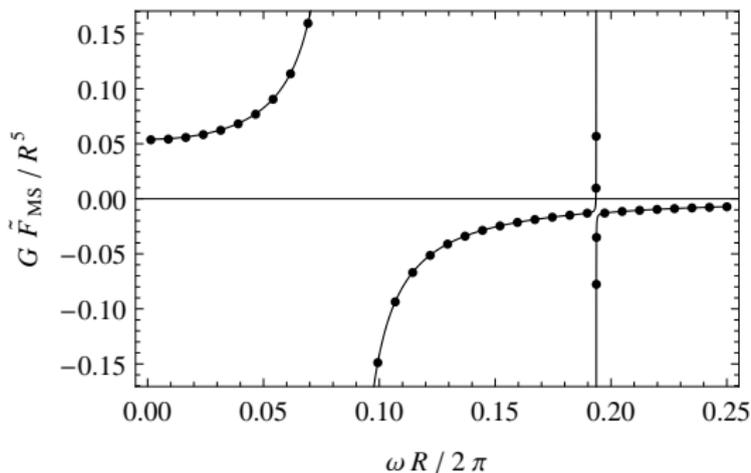
$$X_{IR}^\nu = e^{-i\omega r} (\omega r)^\nu \left(1 - \frac{2M}{r}\right)^{-i2M\omega} \sum_{n=-\infty}^{\infty} \cdots \times (\omega r)^n {}_1F_1(\dots; 2i\omega r)$$

- renormalized angular momentum:  $\nu = \nu(l, M\omega)$  transcendental number

- Nonanalytic terms identical  $\Rightarrow$  **matching**:  $X_{UV}^\nu = K_\nu X_{IR}^\nu$
- Sets on indep. solutions:  $X_{UV}^\nu, X_{UV}^{-\nu-1}$  or  $X_{IR}^\nu, X_{IR}^{-\nu-1}$
- From numerical NS perturbation:  $X = A_1 X_{UV}^\nu + A_2 X_{UV}^{-\nu-1}$
- $X_{IR}^\nu, X_{IR}^{-\nu-1}$  related to effective  $\delta$ -source via **variation of parameters**
- Regularization:  $\delta(\vec{r}) = (rc_l)^{2-l} \frac{\Gamma(\frac{d-\epsilon}{2})}{\pi^{3/2} 2^\epsilon \Gamma(\frac{\epsilon}{2})} \mu_0^\epsilon r^{\epsilon-3}$
- Fit for the response:

$$\tilde{F}_{MS}(\omega) = \sum_n \frac{l_n^2}{\omega_n^2 - \omega^2}$$

- Just like in Newtonian case!
- Relative error less than 2%
- Mode frequencies:  $\omega_n$
- **Relativistic overlap integrals**:  $l_n$
- Matching scale  $\mu_0$  is fitted, too



# Conclusions

## Motivation:

- Adiabatic tidal effects may not be sufficient  
[Maselli, Gualtieri, Pannarale, Ferrari, PRD **86** (2012) 044032]
- Resonances between oscillation modes and orbital motion:
  - **Shattering of NS crust**  
[Tsang, Read, Hinderer, Piro, Bondarescu, PRL **108** (2012) 011102]
  - Numerical simulations of binary NS  
[Gold, Bernuzzi, Thierfelder, Brüggmann, Pretorius, PRD **86** (2012) 121501]
- Definition of source (Dixon) multipoles see also [Harte, CQG **29** (2012) 055012]
- Definition of **relativistic overlap integrals**

## Outlook:

- More realistic NS models: rotation, crust, . . . (also for Newtonian case)
- Dimensional regularization
- Other multipoles based on action in [Goldberger, Ross, PRD **81** (2010) 124015]
- 2nd Love number of **rotating** black holes

Thank you for your attention

and for support by the German Research Foundation **DFG**  
and by the Fundação para a Ciência e a Tecnologia **FCT**